

## THE STRUCTURE OF THE SET OF SINGULAR POINTS OF A CODIMENSION 1 DIFFERENTIAL SYSTEM ON A 5-MANIFOLD

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**ABSTRACT.** Generic modules  $V$  of vector fields tangent to a 5-dimensional smooth manifold  $M$ , generated locally by four not necessarily linearly independent fields  $X_1, X_2, X_3, X_4$ , are considered. Denoting by  $\omega$  the 1-form  $X_4 \lrcorner X_3 \lrcorner X_2 \lrcorner X_1 \lrcorner \overset{5}{\Omega}$  conjugated to  $V$  ( $\overset{5}{\Omega}$  is a fixed local volume form on  $M$ ), the loci of singular behavior of  $V: M_{\deg}(V) = \{p \in M | \omega(p) = 0\}$  and  $M_{\text{sing}}(V) = \{p \in M | \omega \wedge (d\omega)^2(p) = 0\}$  are handled. The local classification of this pair of sets is carried out (outside a curve and a discrete set in  $M_{\deg}(V)$ ) up to a smooth diffeomorphism. In the most complicated case, around points of a codimension 3 submanifold of  $M$ ,  $M_{\text{sing}}(V)$  turns out to be diffeomorphic to the Cartesian product of  $\mathbb{R}^2$  and the Whitney's umbrella in  $\mathbb{R}^3$ .

### 1

We are going to consider generic differential systems of codimension 1 in the tangent bundle over a  $C^\infty$  manifold  $M$  of dimension 5. Unfortunately, this notion in different papers is used in many senses. We mean by it the module generated locally, over the ring of smooth functions on  $M$ , by four vector fields (the fields may happen to be linearly dependent at some points).

The equivalent framework for this investigation is that of singularities of  $k$ -tuples of vector fields on  $\mathbb{R}^n$  set in [JP] (here  $k = 4$ ,  $n = 5$ ). The paper can be considered as a continuation of similar research in dimension 3 (included primarily in [JP]) and in dimension 4 [MR, M1–M3]. The study will be local, so that we shall often use the language of germs of sets, functions, vector fields, etc.

### 2

We shall assume once and for all that all the considered objects are  $C^\infty$  smooth, and this will not be additionally inserted in the statements. A point  $p \in M$  is of interest to us when the germ at  $p$  of the considered system  $V$  is not equivalent to the Darboux model

$$\text{span} \left( \frac{\partial}{\partial y} - u \frac{\partial}{\partial x}, \frac{\partial}{\partial z} - v \frac{\partial}{\partial x}, \frac{\partial}{\partial u}, \frac{\partial}{\partial v} \right),$$

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or, equivalently, to the Pfaffian equation  $dx + u dy + v dz = 0$  (throughout this paper we use  $x, y, z, u, v$  rather than  $x_1, x_2, y_1, y_2, z$ ).

In other words, if (all this locally)  $\Omega$  is a volume form on  $M$  and  $X_1, X_2, X_3, X_4$  generate  $V$ , then one can define a 1-form conjugated to  $V$ ,  $\omega(\cdot) := \Omega(X_1, X_2, X_3, X_4, \cdot)$ , which vanishes at every point  $p$  where  $\dim V(p) \leq 3$ . Such  $\omega$  is not defined invariantly, but the Pfaffian equation it represents already is. So we are interested in either (a)  $\omega(p) = 0$  or (b)  $\omega(p) \neq 0$  and  $\omega \wedge (d\omega)^2|_p = 0$ . The latter means that the class of the Pfaffian equation  $\omega = 0$  at  $p$  is not 5 (regarding this notion, see [F, Ma]).

Throughout this paper the union of the geometric loci of (a) and (b) is denoted by  $M_{\text{sing}}(V)$ , and the locus of (a) alone by  $M_{\text{deg}}(V)$ .

*Note.* We consider only typical degenerations, i.e., those that are unavoidable under arbitrary small perturbations of  $V$ . Therefore we assume that  $\dim V(p) = 3$  for every  $p \in M_{\text{deg}}(V)$  (the falling of  $\dim V(p)$  by 2 is already a codimension 6 feature and is not typical by Transversality theorem, see [AGV, Ma]).

### 3

The normal forms of the smooth classification of germs of generic  $V$ 's were found at generic points of  $M_{\text{sing}}(V) \setminus M_{\text{deg}}(V)$  (i.e., in codimension 1) by Martinet [Ma], and recently at points of certain 2-dimensional surface included in it (where, typically, the next degeneration having codimension 3 materializes) by Zhitomirskii [Z1, Z2]. All these normal forms are simple, i.e., moduleless; such normal forms are called local models.

As regards  $M_{\text{deg}}(V)$ , which typically has codimension 2 (see Proposition in §7), even at its generic points normal forms are unknown. This is in contradistinction to dimensions 3 and 4, where local models at such points of the respective loci of  $\dim V$  falling by 1 were found by Jakubczyk and Przytycki [JP] in the case of  $\dim M = 3$ , and by Mormul and Roussarie [MR] for  $\dim M = 4$ .

We suspect that in dimension 5 at points of  $M_{\text{deg}}(V)$  there are no models, and moreover normal forms contain functional parameters. Yet already the problem of classification of the pair of germs of sets  $M_{\text{sing}}(V)$ ,  $M_{\text{deg}}(V)$  turns out to be interesting. The paper is devoted to this problem.

### 4

**Main Theorem.** *For a generic differential system  $V$  on  $M$  there exist subsets of  $M_{\text{deg}}(V)$ : a curve  $M_1$  and a set of isolated points  $M_2$  such that the germ of  $(M_{\text{sing}}(V), M_{\text{deg}}(V))$  at any point of  $M_{\text{sing}}(V) \setminus (M_1 \cup M_2)$  is equivalent to one of the following germs:*

- (A) *germ of 4-manifold,  $\emptyset$ ;*
- (B) *germ of 3-manifold,  $M_{\text{deg}}(V) = M_{\text{sing}}(V)$ ;*
- (C) *germ of stratified manifold with strata of dimensions 4, 4, 3 (the last in the intersection of closures of the first and second),  $M_{\text{deg}}(V) =$  the 3-dimensional stratum (see Figure 1);*
- (D) *the germ of the Whitney's umbrella  $\times \mathbb{R}^2$ , (its handle)  $\times \mathbb{R}^2$  (see Figure 2).*

*In other words, in suitable coordinates  $x, y, z, u, v$ , the pair of sets  $M_{\text{sing}}(V)$ ,  $M_{\text{deg}}(V)$  is locally given by the equations*

- (A)  $M_{\text{sing}} = \{x = 0\}$ ,  $M_{\text{deg}} = \emptyset$ ;  
 (B)  $M_{\text{sing}} = M_{\text{deg}} = \{x = y = 0\}$ ;  
 (C)  $M_{\text{sing}} = \{xy = 0\}$ ,  $M_{\text{deg}} = \{x = y = 0\}$ ;  
 (D)  $M_{\text{sing}} = \{x^2 = zy^2\}$ ,  $M_{\text{deg}} = \{x = y = 0\}$ .

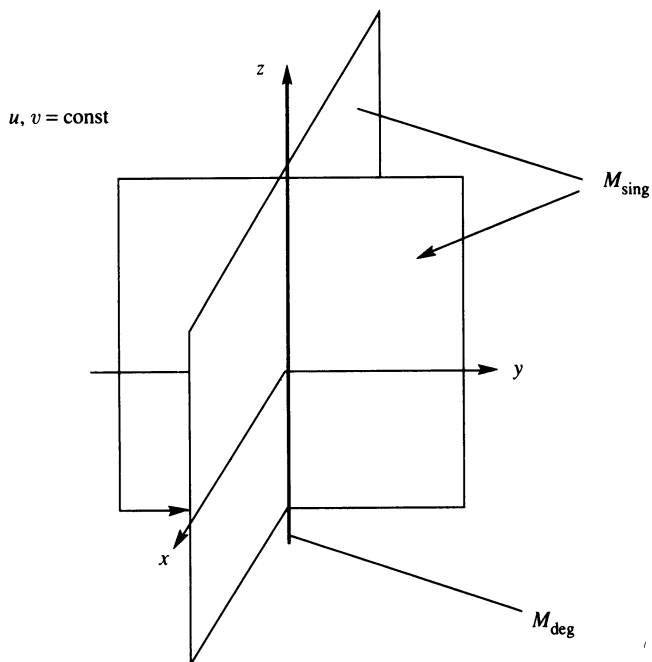


FIGURE 1

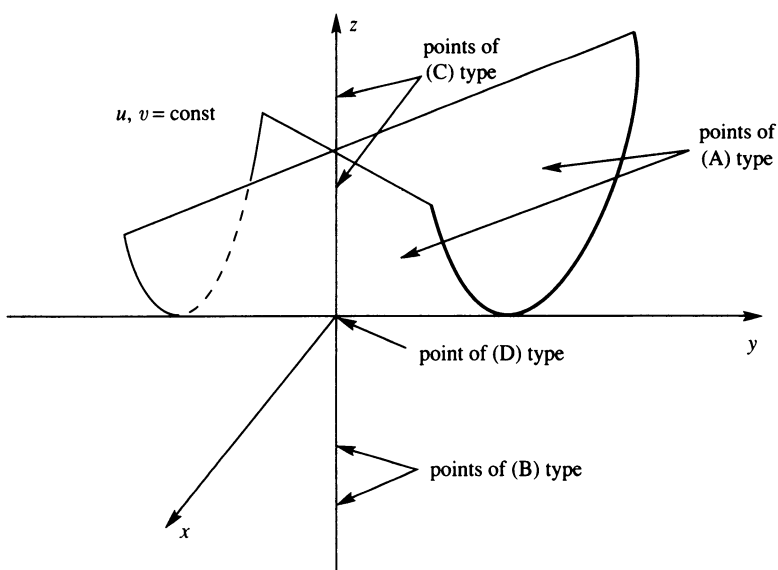


FIGURE 2

In cases (B)–(D), in the mentioned coordinates,

$$V(0) = \text{span}(\partial/\partial x, \partial\partial/\partial y, \partial/\partial z).$$

## 5

As we said in §3, in cases (B)–(D) local models probably do not exist. Nevertheless, we are able then to simplify (locally)  $V$  significantly. This, among other things, will come in handy in proving the Main Theorem in §11.

Let  $\mathfrak{m}_{x,y}$  stand for the ideal of germs at  $0 \in \mathbb{R}^5$  of functions vanishing on  $\{x = y = 0\}$ ;  $\mathfrak{m}_{x,y}^k$  is its  $k$ th power. For brevity we write the same symbol for the set of germs at 0 of 1-forms with coefficients in  $\mathfrak{m}_{x,y}^k$ . In the sequel  $j_{x,y}^k$  will denote the natural projection  $\mathcal{F}_0^5 \rightarrow \mathcal{F}_0^5/\mathfrak{m}_{x,y}^{k+1}$  ( $\mathcal{F}_0^5$  = the ring of germs at 0 of smooth functions on  $\mathbb{R}^5$ ). We shall apply  $j_{x,y}^k$  to germs of 1-forms in the natural sense, too.

The mentioned simplified description of  $V$  (a normal form) is given in terms of the conjugated form  $\omega$  (cf. §2).

**Theorem on normal form.** *Let  $p \in M_{\deg}(V) \setminus M_1$ . There exist coordinates  $x, y, z, u, v$  (vanishing at  $p$ ) s.t. the germ of  $\omega$  at  $p$  has the form*

$$(1) \quad x du + y dv + f dz, \quad f \in \mathfrak{m}_{x,y}^2.$$

Observe that in these coordinates

$$(2) \quad M_{\deg}(V) = \{x = y = 0\}.$$

**Corollary 1.** *At any  $p \in M_{\deg}(V)$ ,  $j_p^1(\omega \wedge (d\omega)^2) = 0$ .*

Indeed, in the normal form coordinates

$$(3) \quad \omega \wedge (d\omega)^2 = 2(xf_x + yf_y - f) dx \wedge dy \wedge dz \wedge du \wedge dv,$$

and  $f, xf_x, yf_y \in \mathfrak{m}_{x,y}^2$ . (Here and in the sequel the symbol of a function followed by a lowercase letter subscript denotes the respective function's partial derivative.)  $\square$

**Remark 1.** In the normal form coordinates Corollary 1 can be written compactly as  $j_{x,y}^1(\omega \wedge (d\omega)^2) = 0$ .

**Corollary 2.** *Let  $p \in M_{\deg}(V) \setminus M_1$ . There exist coordinates  $x, y, z, u, v$  such that the germ of  $V$  at  $p$  is generated by vector fields  $\partial/\partial x, \partial/\partial y, \partial/\partial z + f_1\partial/\partial u + f_2\partial/\partial v, y\partial/\partial u - x\partial/\partial v$ , where  $f_1, f_2 \in \mathfrak{m}_{x,y}$ .*

*Proof.* Let  $\omega$  be the 1-form conjugated to  $V$  and  $x, y, z, u, v$  be the coordinates of n.f. (1). We can write  $f = -xf_1 - yf_2$  with  $f_1, f_2 \in \mathfrak{m}_{x,y}$ . Any v.f.  $\xi$  from  $V$ ,

$$\xi = A_1 \frac{\partial}{\partial x} + A_2 \frac{\partial}{\partial y} + A_3 \frac{\partial}{\partial z} + A_4 \frac{\partial}{\partial u} + A_5 \frac{\partial}{\partial v},$$

satisfies  $\omega(\xi) \equiv 0$  or, equivalently,  $xA_4 + yA_5 - (xf_1 + yf_2)A_3 \equiv 0$ . The latter relation implies  $A_4 - f_1A_3 = yg$  and  $A_5 - f_2A_3 = -xg$  for some function  $g$ . Thus

$$\xi = A_1 \frac{\partial}{\partial x} + A_2 \frac{\partial}{\partial y} + A_3 \left( \frac{\partial}{\partial z} + f_1 \frac{\partial}{\partial u} + f_2 \frac{\partial}{\partial v} \right) + g \left( y \frac{\partial}{\partial u} - x \frac{\partial}{\partial v} \right). \quad \square$$

**Remark 2.** It follows from Corollary 2 that the generators of  $V$  and their first order Lie brackets span at 0 the full 5-dimensional tangent space. (This prop-

erty of  $V$  is weaker than the transversality assumed in §8 and occurs also without that transversality.)

We postpone the proof of the theorem on normal form till §12; the final part of it occupies §13.

## 6

In proving the Main Theorem we shall represent  $V$  by a triple  $(\omega_1, \omega_2, Y)$ , where  $\omega_1, \omega_2$  are 1-forms, and  $Y$  is a vector field.

**Observation 1.** *For  $p \in M_{\deg}(V)$  there exist germs at  $p$  of a vector field  $Y$  and of independent Pfaffian forms  $\omega_1$  and  $\omega_2$  such that*

$$(4) \quad \omega := Y \lrcorner (\omega_1 \wedge \omega_2)$$

*is conjugated to  $V$  and  $M_{\deg}(V) = \{\omega_1(Y) = \omega_2(Y) = 0\}$ .*

Passing to the  $(\omega_1, \omega_2, Y)$ 's is purposeful, since any  $\omega$  conjugated to  $V$  is highly nontypical among all differential 1-forms (it vanishes on a "big" set  $M_{\deg}(V)$ ), while the objects in the triple  $(\omega_1, \omega_2, Y)$  do not vanish at any point.

*Proof.* Let  $V$  be generated by four vector fields  $X_1, X_2, X_3, Y$ , three of which are independent (for instance,  $X_1, X_2, X_3$ ). The distribution spanned by them can be described by a Pfaffian system  $\omega_1 = \omega_2 = 0$ , and then  $V = \text{span}(Y, \ker \omega_1 \cap \ker \omega_2)$ . Now the statements of the observation are easy to verify.  $\square$

It seems to us that the above representation of  $\omega$  is of its own interest and can be applied in many situations.

## 7

Our first step in the proof of the Main Theorem consists in ensuring that

**Proposition.** *For  $V$  generic  $M_{\deg}(V)$  is, if not empty, a smooth codimension 2 submanifold of  $M$ .*

*Proof.* We use description (4) and consider the set  $Q_1$  of 1-jets at  $0 \in \mathbb{R}^5$  of 3-tuples  $(\omega_1, \omega_2, Y)$ . Let its subset  $\tilde{Q}_1$  be given by equations

$$\omega_1(Y)(0) = \omega_2(Y)(0) = 0, \quad d(\omega_1(Y)) \wedge d(\omega_2(Y))|_0 = 0.$$

Clearly  $\tilde{Q}_1$  has codimension 6 in  $Q_1$ . The standard use of Transversality theorem gives that for generic  $V$  its 1-jet is nowhere included in  $\tilde{Q}_1$ . This implies the conclusion of the Proposition.  $\square$

*Remark 3.* The theory mentioned in §1, developed in [JP], yields, among many other things, that the locus of  $\dim V = k - 1$  is smooth for a generic codimension  $n - k$  differential system in  $\mathbb{R}^n$ .

## 8

Here is the definition (in invariant terms) of the sets  $M_1, M_2 \subset M_{\deg}(V)$  occurring in the Main Theorem.

A point  $p$  is included in  $M_1$  iff  $V(p)$  is not transversal to  $M_{\deg}(V)$ ; a point  $p$  is included in  $M_2$  iff  $j_p^2(\omega \wedge (d\omega)^2) = 0$ .

Now we are going to prove

(i)  $M_1$  is a smooth curve.

The proof is based on the following

**Lemma 1.** *For a generic system  $V$ , at any  $p \in M_{\deg}(V)$ ,  $V(p)$  is transversal to  $M_{\deg}(V) \Leftrightarrow \text{rank}(d\omega)^2(p) = 4$ .*

*Proof.* We take the  $\omega$  from Observation 1. Computing at a point of  $M_{\deg}(V)$  and using Observation 1, we get  $d\omega = d(\omega_1(Y)) \wedge \omega_2 - d(\omega_2(Y)) \wedge \omega_1$ . Thus  $(d\omega)^2 = -2\omega_1 \wedge \omega_2 \wedge d(\omega_1(Y)) \wedge d(\omega_2(Y))$ . The condition  $\text{rank}(d\omega)^2 = 4$  means that the kernels of the 1-forms entering the above formula intersect one another as sparingly as possible.

Because for  $p \in M_{\deg}(V)$ ,  $V(p) = \ker \omega_1(p) \cap \ker \omega_2(p)$ , and since, in view of  $d(\omega_1(Y)) \wedge d(\omega_2(Y)) \neq 0$ ,  $M_{\deg}(V)$  is a smooth manifold and  $T_p M_{\deg}(V) = \ker d(\omega_1(Y)) \cap \ker d(\omega_2(Y))$  (cf. Proposition in §7 and Observation 1), we conclude that so intersect each other  $V(p)$  and  $T_p M_{\deg}(V)$ , which means transversality.

The implication  $\Rightarrow$  uses the same arguments and the condition of smoothness of  $M_{\deg}(V)$  valid for a generic system  $V$ .  $\square$

*Note.* The statement on the right-hand side of Lemma 1 is here equivalent to  $(d\omega)^2(p) \neq 0$ .

*Proof of (i).* Let  $p \in M_1$ . By Lemma 1  $\text{rank}(d\omega)^2(p) < 4$ , which means that the 1-jet at  $p$  of  $(\omega_1, \omega_2, Y)$  satisfies the conditions

$$\omega_1(Y)(p) = 0, \quad \omega_2(Y)(p) = 0, \quad \omega_1 \wedge \omega_2 \wedge d(\omega_1(Y)) \wedge d(\omega_2(Y))|_p = 0.$$

It is clear that these conditions distinguish a (stratified) codimension 4 manifold in the space  $Q_1$  of 1-jets. One can show that its singular points form a set of codimension 6 in  $Q_1$ . Therefore (i) follows from Transversality theorem.  $\square$

In turn, we can show that

(ii)  $M_2$  consists of isolated points.

*Proof.*  $M_2$  is defined invariantly (see §8), and we prefer to work in the coordinates giving (1). Let

$$(5) \quad f = A(z, u, v)x^2 + B(z, u, v)xy + C(z, u, v)y^2 + \tilde{f}, \quad \tilde{f} \in \mathfrak{m}_{x,y}^3.$$

Then in view of (2) and (3)  $M_2$  is given locally by the equations

$$x = y = A = B = C = 0.$$

The application of Transversality theorem completes the proof.  $\square$

## 11. PROOF OF THE MAIN THEOREM

(A) This is the well-known case: for  $V$  generic in the sense of being transversal (after natural identifications) to the stratification  $C$  of 1-jets of the Pfaffian equations on  $M$ , constructed in [Ma, p. 136],  $M_{\text{sing}}(V) \setminus M_{\deg}(V)$  is a smooth codimension 1 submanifold of  $M$ .

Now consider the case  $p \in M_{\deg}(V) \setminus (M_1 \cup M_2)$ . We are using the normal form (1) coordinates. Putting  $g := xf_x + yf_y - f$  and using (5), one has

$$g = A(z, u, v)x^2 + B(z, u, v)xy + C(z, u, v)y^2 + \hat{f}, \quad \hat{f} \in \mathfrak{m}_{x,y}^3.$$

By (3),  $(M_{\text{sing}}(V), M_{\deg}(V)) = (\{g = 0\}, \{x = y = 0\})$ . Now to prove the Main Theorem it suffices to prove that there exists a local diffeomorphism  $\Phi$  preserving the manifold  $\{x = y = 0\}$  and such that  $g \circ \Phi = \mp(x^2 + y^2)$ , or  $g \circ \Phi = xy$ , or  $g \circ \Phi = x^2 - y^2z$ . As we are outside  $M_2$ ,  $j_p^2 g \neq 0$ . We can assume  $p = 0 \in \mathbb{R}^5$ . Suppose at first that

$$(6) \quad \begin{vmatrix} 2A & B \\ B & 2C \end{vmatrix} (0) \neq 0.$$

Two subcases are possible:

(B) the Hessian is positively or negatively defined, and

(C) (6) holds and the Hessian is neither positively nor negatively defined.

By the Morse lemma with parameters [AGV] we can claim the existence of a coordinate change

$$(7) \quad x \rightarrow x + \phi(x, y, z), \quad y \rightarrow y + \psi(x, y, z), \quad \phi, \psi \in \mathfrak{m}_{x,y}^3,$$

simplifying the function  $g$  to  $x^2 + y^2$  or  $-x^2 - y^2$  in case (B) and to  $xy$  in (C). Such transformation (7) preserves (2), and one obtains normal forms (B) and (C) in §4.

Now suppose that condition (6) is violated. Since  $j^2 g \neq 0$ , we can assume that  $A(0) \neq 0$ , and further that  $A(z, u, v) \equiv 1$ . Using again the Morse lemma with parameters (the parameters are now  $y, z, u, v$ ) we can claim the existence of a coordinate change of form (7) bringing the function  $g$  to the form

$$(8) \quad g = x^2 + y^2\tau(z, u, v) + y^3\nu(y, z, u, v), \quad \tau(0) = 0.$$

For a generic differential system  $d\tau(0) \neq 0$  (the condition  $d\tau(0) = 0$  together with violating (6) and the inclusion of the source point 0 in  $M_{\deg}(V)$  give the degeneration of codimension 6, nontypical by Transversality theorem). This condition is a bit stronger than the following one: the set of points  $p \in M_{\deg}(V)$  such that the germ at  $p$  of a pair  $(M_{\text{sing}}(V), M_{\deg}(V))$  is not equivalent either to normal form (B) or (C) in §4 is a smooth 2-dimensional submanifold. Therefore there exists a transformation of the form

$$(9) \quad z \rightarrow \eta_1(z, u, v), \quad u \rightarrow \eta_2(z, u, v), \quad v \rightarrow \eta_3(z, u, v)$$

simplifying expression (8) of  $g$  to

$$(10) \quad g = x^2 - y^2z + y^3\tilde{\nu}(y, z, u, v).$$

A transformation of the form

$$z \rightarrow z + y\alpha(y, z, u, v)$$

reduces  $\tilde{\nu}$  in (10) to

$$\hat{\nu} = \tilde{\nu}(y, z + y\alpha, u, v) - \alpha.$$

By the implicit function theorem we can ensure  $\hat{\nu} = 0$  by choosing a suitable function  $\alpha$ . Thus we arrive at normal form (D) in §4.

Finally, it is clear that normalizing the equation for  $M_{\text{sing}}(V)$  above we preserve the description of  $V(0) = \text{span}(\partial/\partial x, \partial/\partial y, \partial/\partial z)$  valid for system  $V$  in normal form (cf. Corollary 2). The proof of the Main Theorem is complete.  $\square$

## 12

The proof of the theorem on normal form will be split into several assertions in this section and the next.

**Lemma 2.** *If  $(d\omega)^2(0) \neq 0$  then  $\omega$  is reducible to  $x du + y dv + \tau$ , where  $\tau \in \mathfrak{m}_{x,y}^2$ .*

*Proof.* Let us choose such coordinates that (2) holds. Then obviously  $j_{x,y}^0 \omega = 0$  and one has an expansion

$$(11) \quad j_{x,y}^1 \omega = (xA_{11} + yA_{12})dx + (xA_{21} + yA_{22})dy + (xA_{31} + yA_{32})dz \\ + (xA_{41} + yA_{42})du + (xA_{51} + yA_{52})dv,$$

where  $A_{ij} = A_{ij}(z, u, v)$ . Now we consider two 1-forms on  $\mathbb{R}^3(z, u, v)$ :

$$\mu_i := A_{3i}(z, u, v)dz + A_{4i}(z, u, v)du + A_{5i}(z, u, v)dv, \quad i = 1, 2.$$

The transformation

$$(12) \quad \begin{aligned} x &\rightarrow \alpha_{11}(z, u, v) \cdot x + \alpha_{12}(z, u, v) \cdot y, \\ y &\rightarrow \alpha_{21}(z, u, v) \cdot x + \alpha_{22}(z, u, v) \cdot y \end{aligned}$$

brings  $j_{x,y}^1 \omega$  to the form

$$(\tilde{A}_{11}x + \tilde{A}_{12}y)dx + (\tilde{A}_{21}x + \tilde{A}_{22}y)dy + x\tilde{\mu}_1 + y\tilde{\mu}_2,$$

where

$$\tilde{\mu}_1 = \alpha_{11}\mu_1 + \alpha_{12}\mu_2, \quad \tilde{\mu}_2 = \alpha_{21}\mu_1 + \alpha_{22}\mu_2.$$

Direct calculation shows that the assumption  $(d\omega)^2(0) \neq 0$  means exactly  $\mu_1 \wedge \mu_2|_0 \neq 0$  (the reason we have introduced  $\mu_1, \mu_2$ !). Apply any transformation (9) that straightens in  $\mathbb{R}^3$  the line field  $\ker \mu_1 \wedge \mu_2$  to  $\text{span}(\partial/\partial z)$ . Understanding it as the coordinate change in  $\mathbb{R}^5$ , one has then  $\text{span}(\mu_1, \mu_2) = \text{span}(du, dv)$  in a neighbourhood of  $0 \in \mathbb{R}^5$ . This can be improved, using (12), to  $\tilde{\mu}_1 = du, \tilde{\mu}_2 = dv$ , yielding

$$j_{x,y}^1 \omega = (x\tilde{A}_{11} + y\tilde{A}_{12})dx + (x\tilde{A}_{21} + y\tilde{A}_{22})dy + xdu + ydv.$$

The change of coordinates

$$u \rightarrow u - x\tilde{A}_{11} - y\tilde{A}_{21}, \quad v \rightarrow v - x\tilde{A}_{12} - y\tilde{A}_{22}$$

eventually annihilates  $j_{x,y}^1(\omega - xdu - ydv)$ .  $\square$

Lemma 2 will serve as the premise for  $k = 1$  in the inductive argument justifying Corollary 3 (see below). The following will constitute the induction step.



**Lemma 3.** For any  $k \geq 1$  and  $\omega$  satisfying  $(d\omega)^2(0) \neq 0$  the jet  $j_{x,y}^k \omega$  is reducible to form (1).

*Proof.* Suppose that for certain  $k \geq 1$   $j_{x,y}^k \omega$  is already reduced to form (1). We shall show that  $j_{x,y}^{k+1} \omega$  can be so reduced, too. Let  $\mathfrak{m}_{x,y}^{(l)}$  be the set of all function germs of the form  $\sum_{i+j=l} c_{ij}(z, u, v)x^i y^j$ , where  $c_{ij}$  are germs at  $0 \in \mathbb{R}^3$  of smooth functions. Write

$$(13) \quad j_{x,y}^{k+1} \omega = j_{x,y}^k \omega + A_1 dx + A_2 dy + A_3 dz + A_4 du + A_5 dv,$$

where  $A_i \in \mathfrak{m}_{x,y}^{(k+1)}$ . We are going to take new coordinates

$$(14) \quad x + \varphi, \quad y + \psi, \quad z, \quad u + \gamma, \quad v + \mu, \quad \text{where } \varphi, \psi, \gamma, \mu \in \mathfrak{m}_{x,y}^{(k+1)}.$$

Denoting by  $T$  the right-hand side of (13),  $j_{x,y}^{k+1} \omega$  assumes in these coordinates the form  $T + (x\gamma_x + y\mu_x)dx + (x\gamma_y + y\mu_y)dy + \varphi du + \psi dv$ . So it suffices to solve the system

$$(15) \quad \varphi + A_3 = \psi + A_4 = 0, \quad x\gamma_x + y\mu_x + A_1 = 0, \quad x\gamma_y + y\mu_y + A_2 = 0.$$

Obviously  $\varphi = -A_3$ ,  $\psi = -A_4$ . By putting  $R := x\gamma + y\mu$  we reduce (15) to the system of three equations for  $R$ ,  $\gamma$ ,  $\mu$ :

$$R_x - \gamma + A_1 = 0, \quad R_y - \mu + A_2 = 0, \quad R = x\gamma + y\mu,$$

$\gamma, \mu \in \mathfrak{m}_{x,y}^{(k+1)}$ . This can be written briefly as

$$(16) \quad R = x(R_x + A_1) + y(R_y + A_2), \quad R \in \mathfrak{m}_{x,y}^{(k+2)}$$

(having such  $R$ , we take  $\gamma := R_x + A_1$ ,  $\mu := R_y + A_2$ ). As  $x A_1 + y A_2 = \sum_{i+j=k+2} b_{ij}(z, u, v)x^i y^j$ , we can give an explicit solution to (16):

$$R = \sum_{i+j=k+2} (1-i-j)^{-1} b_{ij}(z, u, v)x^i y^j. \quad \square$$

We denote by  $\mathfrak{m}_{x,y}^\infty$  the ideal of germs of functions vanishing on  $\{x = y = 0\}$  together with all their partial derivatives, and also the set of 1-form germs having such coefficients, and the set of respective vector field germs, too. Consequently,  $j_{x,y}^\infty$  is defined analogously to  $j_{x,y}^k$  (see §5).

**Corollary 3.**  $\omega$  as in Lemma 3 is reducible to  $x du + y dv + f dz + \tau$ , where  $\tau$  is a 1-form,  $\tau \in \mathfrak{m}_{x,y}^\infty$ , and  $f \in \mathfrak{m}_{x,y}^2$ .

*Proof.* Using Lemma 2 as the departure point ( $k = 1$ ) we reduce inductively consecutive jets  $j_{x,y}^k \omega$ , applying Lemma 3 at each step. Taking into account the character of normalizing transformations (14) and the fact that, after passing to a fixed representative of  $\omega$ , (16) is solvable in an independent of  $k$  neighbourhood of  $0 \in \mathbb{R}^3(z, u, v)$ , there exists a formal in  $x, y$  transformation of  $\mathbb{R}^5$ , having as coefficients (of its series in  $x, y$ ) smooth functions of  $z, u, v$  defined in a common neighbourhood of  $0$ , which reduces  $j_{x,y}^\infty \omega$  to form (1). Now it suffices to apply the Whitney extension theorem (see [W]).  $\square$

To prove the theorem on normal form we must still prove

**Lemma 4.** *Let  $f \in \mathfrak{m}_{x,y}^2$ ,  $\tau$  be a 1-form,  $\tau \in \mathfrak{m}_{x,y}^\infty$ . Then the 1-form  $\omega = x du + y dv + f dz + \tau$  is reducible to the form  $x du + y dv + \tilde{f} dz$ ,  $\tilde{f} \in \mathfrak{m}_{x,y}^2$ .*

*Proof.* We use some modifications of the homotopy method [Z2, Chapter 1, §3]. Let  $\hat{\omega} := x du + y dv$ . Introduce also the truncation operator  $P$  sending every 1-form  $\kappa_1 dx + \kappa_2 dy + \kappa_3 dz + \kappa_4 du + \kappa_5 dv$  ( $\kappa_i$  are functions of  $x, y, z, u, v$ ) to  $\kappa_1 dx + \kappa_2 dy + \kappa_4 du + \kappa_5 dv$ , and the family of forms  $\omega_t := \hat{\omega} + t(f dz + \tau)$ ,  $t \in [0, 1]$ . Consider the equation

$$(17) \quad P(X_t \lrcorner d\omega_t + d(X_t \lrcorner \omega_t) + f dz + \tau) = 0$$

for an unknown family of vector fields  $X_t$ . (We shall require additionally that  $X_t \lrcorner dz \equiv 0$ .) Equation (17) can be equivalently written as

$$(18) \quad P(X_t \lrcorner d\omega_t + d(X_t \lrcorner \omega_t) + \tau) = 0.$$

*Claim.* If there exists a smooth family  $X_t \in \mathfrak{m}_{x,y}^\infty$  depending on  $t$  and satisfying (18) and such that  $X_t \lrcorner dz \equiv 0$ , then Lemma 4 holds. (Compare the classical variant of the homotopy method [AGV].) In order to substantiate the Claim, consider the family of diffeomorphisms  $\phi_t$  defined by

$$\frac{d\phi_t}{dt} = X_t(\phi_t), \quad \phi_0 = \text{id}.$$

Then

$$\begin{aligned} \frac{d}{dt}(P(\phi_t^* \omega_t)) &= P\left(\frac{d}{dt}(\phi_t^* \omega_t)\right) = P\phi_t^* \left(L_{X_t} \omega_t + \frac{d\omega_t}{dt}\right) \\ &= P\phi_t^* (X_t \lrcorner d\omega_t + d(X_t \lrcorner \omega_t) + f dz + \tau), \end{aligned}$$

where  $L_{X_t} \omega_t$  is the Lie derivative of  $\omega_t$  along the field  $X_t$ . Equation (17) implies that  $L_{X_t} \omega_t + d\omega_t/dt \in \ker P \ \forall t$ . By virtue of  $X_t \lrcorner dz \equiv 0$ , also  $\phi_t^*(L_{X_t} \omega_t + d\omega_t/dt)$  is always included in  $\ker P$ ; hence  $d(P(\phi_t^* \omega_t))/dt \equiv 0$ . This infers  $P\phi_1^* \omega_1 = P\phi_0^* \omega_0$ , or  $\phi_1^* \omega - \hat{\omega} \in \ker P$ , i.e.,  $\phi_1^* \omega - \hat{\omega} = \tilde{f} dz$ . Noticing that  $\phi_t = \text{id} + \varphi_t$ ,  $\varphi_t \in \mathfrak{m}_{x,y}^\infty$ , the assumption  $f \in \mathfrak{m}_{x,y}^2$  obviously implies  $\tilde{f} \in \mathfrak{m}_{x,y}^2$ , proving the Claim.

Now we are going to show that the premise in the Claim holds. Seeking  $X_t$  in the form  $f_{1,t} \partial/\partial x + f_{2,t} \partial/\partial y + f_{3,t} \partial/\partial u + f_{4,t} \partial/\partial v$ ,  $f_{i,t} \in \mathfrak{m}_{x,y}^\infty$ , (18) boils down to a system of five equations for the  $f_{i,t}$ 's and  $R_t := X_t \lrcorner \omega_t$ , the last defining identity being itself the fifth equation which, on writing  $\tau = \tau_1 dx + \tau_2 dy + \tau_3 du + \tau_4 dv$  and computing  $X_t \lrcorner \omega_t$  explicitly, assumes the form

$$R_t = t\tau_1 f_{1,t} + t\tau_2 f_{2,t} + (x + t\tau_3) f_{3,t} + (y + t\tau_4) f_{4,t}.$$

The unknowns  $f_{i,t}$  can be eliminated from the first four equations (they can be expressed via the first order partial derivatives of  $R_t$ ), after which we arrive at one equation for  $R_t$ :

$$(19) \quad R_t - x(R_t)_x - y(R_t)_y + \Theta_t(R_t) = a_t,$$

where  $\Theta_t$  is a family of vector fields,  $a_t$  is a family of function germs,  $\Theta_t \in \mathfrak{m}_{x,y}^\infty$ ,  $a_t \in \mathfrak{m}_{x,y}^\infty$ . (Observe that (19) is, to some extent, similar to (16), but the occurring flat function and vector field make the great difference.)

The fields  $-x\partial/\partial x - y\partial/\partial y + \Theta_t$  are hyperbolic on the manifold  $\{x = y = 0\}$  (this manifold is attracting for the respective dynamical system). Thanks to that, by virtue of the Belitskii results (see [B1, B2] and references in the latter<sup>1</sup>), (19) has a smooth family of solutions  $R_t \in \mathfrak{m}_{x,y}^\infty$  depending on  $t$ . The proof of Lemma 4 is finished.  $\square$

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<sup>1</sup>Various results by Belitskii on the solvability of singular partial differential equations can be found in [Z2, Chapter 1, §6], where they are collected in a compact form.